# Compressible Flow - TME085 Lecture 10

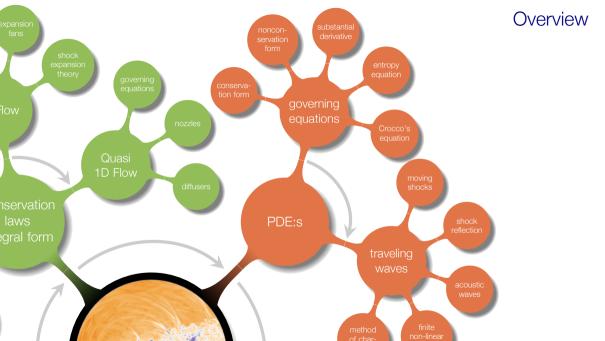
#### Niklas Andersson

Chalmers University of Technology
Department of Mechanics and Maritime Sciences
Division of Fluid Mechanics
Gothenburg, Sweden

niklas.andersson@chalmers.se



# Chapter 7 Unsteady Wave Motion

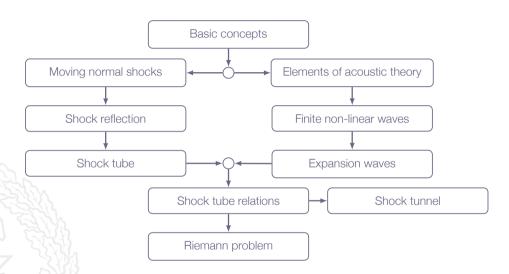


### **Learning Outcomes**

- 3 Describe typical engineering flow situations in which compressibility effects are more or less predominant (e.g. Mach number regimes for steady-state flows)
- 4 Present at least two different formulations of the governing equations for compressible flows and explain what basic conservation principles they are based on
- 8 Derive (marked) and apply (all) of the presented mathematical formulae for classical gas dynamics
  - a 1D isentropic flow\*
  - b normal shocks\*
  - unsteady waves and discontinuities in 1D
  - k basic acoustics
  - Solve engineering problems involving the above-mentioned phenomena (8a-8k)
- Explain how the equations for aero-acoustics and classical acoustics are derived as limiting cases of the compressible flow equations

moving normal shocks - frame of reference seems to be the key here?!

### Roadmap - Unsteady Wave Motion



#### Motivation

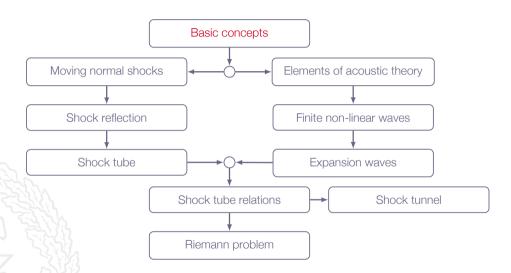
Most practical flows are unsteady

Traveling waves appears in many real-life situations and is an important topic within compressible flows

We will study unsteady flows in one dimension in order to reduce complexity and focus on the physical effects introduced by the unsteadiness

Throughout this section, we will study an application called the shock tube, which is a rather rare application but it lets us study unsteady waves in one dimension and it includes all physical principles introduced in chapter 7

### Roadmap - Unsteady Wave Motion



# **Unsteady Wave Motion**

#### inertial frames!

Physical laws are the same for both frame of references

Shock characteristics are the same for both observers (shape, strength, etc)

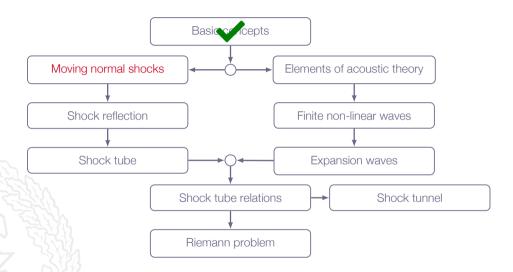
# **Unsteady Wave Motion**

Is there a connection with stationary shock waves?

Answer: Yes!

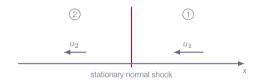
Locally, in a moving frame of reference, the shock may be viewed as a stationary normal shock

### Roadmap - Unsteady Wave Motion



# Chapter 7.2 Moving Normal Shock Waves

Chapter 3: stationary normal shock



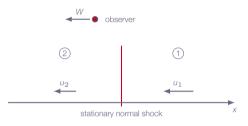
 $u_1 > a_1$  (supersonic flow)

 $u_2 < a_2$  (subsonic flow)

 $p_2 > p_1$  (sudden compression)

 $s_2 > s_1$  (shock loss)





- ► Introduce observer moving to the left with speed W
  - ▶ if *W* is constant the observer is still in an inertial system
  - all physical laws are unchanged
- The observer sees a normal shock moving to the right with speed W
  - $\triangleright$  gas velocity ahead of shock:  $u'_1 = W u_1$
  - ightharpoonup gas velocity behind shock:  $u_2' = W u_2$

Now, let 
$$W = u_1 \Rightarrow$$

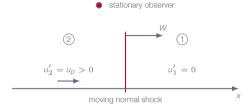
$$u_1' = 0$$

$$u_2' = u_1 - u_2 > 0$$

The observer now sees the shock traveling to the right with speed  $W = u_1$  into a stagnant gas, leaving a compressed gas  $(p_2 > p_1)$  with velocity  $u_2' > 0$  behind it

Introducing  $u_p$ :

$$u_p = u_2' = u_1 - u_2$$



#### Case 1

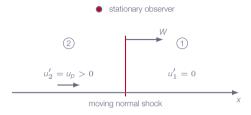
Analogy:

- stationary normal shock
- observer moving with velocity W

#### Case 2

- normal shock moving with velocity W
- stationary observer

# Moving Normal Shock Waves - Governing Equations



For stationary normal shocks we have:

$$\rho_1 u_1 = \rho_2 u_2$$

$$\rho_1 u_1^2 + \rho_1 = \rho_2 u_2^2 + \rho_2$$

$$h_1 + \frac{1}{2} u_1^2 = h_2 + \frac{1}{2} u_2^2$$

With  $(u_1 = W)$  and  $(u_2 = W - u_p)$  we get:

$$\rho_1 W = \rho_2 (W - u_p)$$

$$\rho_1 W^2 + \rho_1 = \rho_2 (W - u_p)^2 + \rho_2$$

$$h_1 + \frac{1}{2} W^2 = h_2 + \frac{1}{2} (W - u_p)^2$$

Starting from the governing equations

$$\rho_1 W = \rho_2 (W - u_\rho)$$

$$\rho_1 W^2 + \rho_1 = \rho_2 (W - u_\rho)^2 + \rho_2$$

$$h_1 + \frac{1}{2} W^2 = h_2 + \frac{1}{2} (W - u_\rho)^2$$

and using 
$$h = e + \frac{p}{\rho}$$

it is possible to show that

$$e_2 - e_1 = \frac{\rho_1 + \rho_2}{2} \left( \frac{1}{\rho_1} + \frac{1}{\rho_2} \right)$$

$$e_2 - e_1 = \frac{\rho_1 + \rho_2}{2} \left( \frac{1}{\rho_1} + \frac{1}{\rho_2} \right)$$

same Hugoniot equation as for stationary normal shock

This means that we will have same shock strength, *i.e.* same jumps in density, velocity, pressure, etc

Starting from the Hugoniot equation one can show that

$$\frac{\rho_2}{\rho_1} = \frac{1 + \frac{\gamma + 1}{\gamma - 1} \left(\frac{\rho_2}{\rho_1}\right)}{\frac{\gamma + 1}{\gamma - 1} + \frac{\rho_2}{\rho_1}}$$

and

$$\frac{T_2}{T_1} = \frac{\rho_2}{\rho_1} \left[ \frac{\frac{\gamma+1}{\gamma-1} + \frac{\rho_2}{\rho_1}}{1 + \frac{\gamma+1}{\gamma-1} \left(\frac{\rho_2}{\rho_1}\right)} \right]$$

For calorically perfect gas and stationary normal shock:

$$\frac{\rho_2}{\rho_1} = 1 + \frac{2\gamma}{\gamma + 1} (M_s^2 - 1)$$

same as eq. (3.57) in Anderson with  $M_1 = M_S$ 

where

$$M_{\rm S} = \frac{W}{a_1}$$

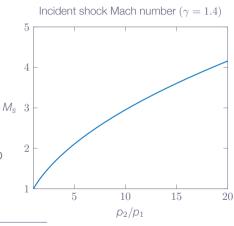
- $ightharpoonup M_s$  is simply the speed of the shock (W), traveling into the stagnant gas, normalized by the speed of sound in this stagnant gas ( $a_1$ )
  - $M_s > 1$ , otherwise there is no shock!
  - shocks always moves faster than sound no warning before it hits you 😊

$$\frac{\rho_2}{\rho_1} = 1 + \frac{2\gamma}{\gamma + 1} (M_s^2 - 1)$$

Re-arrange ⇒

$$M_{\rm S} = \sqrt{\frac{\gamma + 1}{2\gamma} \left(\frac{\rho_2}{\rho_1} - 1\right) + 1}$$

shock speed directly linked to pressure ratio



$$M_{S} = \frac{W}{a_1} \Rightarrow W = a_1 M_{S} = a_1 \sqrt{\frac{\gamma + 1}{2\gamma} \left(\frac{\rho_2}{\rho_1} - 1\right) + 1}$$

From the continuity equation we get:

$$u_{p} = W\left(1 - \frac{\rho_{1}}{\rho_{2}}\right) > 0$$

After some derivation we obtain:

$$u_p = \frac{a_1}{\gamma} \left( \frac{\rho_2}{\rho_1} - 1 \right) \left[ \frac{\frac{2\gamma}{\gamma + 1}}{\frac{\rho_2}{\rho_1} + \frac{\gamma - 1}{\gamma + 1}} \right]^{1/2}$$

Induced Mach number:

$$M_{p} = \frac{u_{p}}{a_{2}} = \frac{u_{p}}{a_{1}} \frac{a_{1}}{a_{2}} = \frac{u_{p}}{a_{1}} \sqrt{\frac{T_{1}}{T_{2}}}$$

inserting  $u_0/a_1$  and  $T_1/T_2$  from relations on previous slides we get:

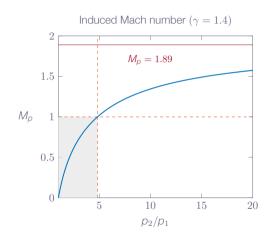
$$M_{\rho} = \frac{1}{\gamma} \left( \frac{\rho_2}{\rho_1} - 1 \right) \left[ \frac{\frac{2\gamma}{\gamma + 1}}{\frac{\gamma - 1}{\gamma + 1} + \frac{\rho_2}{\rho_1}} \right]^{1/2} \left[ \frac{1 + \left( \frac{\gamma + 1}{\gamma - 1} \right) \left( \frac{\rho_2}{\rho_1} \right)}{\left( \frac{\gamma + 1}{\gamma - 1} \right) \left( \frac{\rho_2}{\rho_1} \right) + \left( \frac{\rho_2}{\rho_1} \right)^2} \right]^{1/2}$$

#### Note!

$$\lim_{\frac{\rho_2}{\rho_1} \to \infty} M_{\rho} \to \sqrt{\frac{2}{\gamma(\gamma - 1)}}$$

for air ( $\gamma = 1.4$ )

 $\lim_{\frac{\rho_2}{\to \infty}} M_{\rho} \to 1.89$ 



Moving normal shock with  $p_2/p_1 = 10$ 

$$(p_1 = 1.0 \text{ bar}, T_1 = 300 \text{ K}, \gamma = 1.4)$$

$$\Rightarrow$$
  $M_{\rm S}=2.95$  and  $W=1024.2~m/s$ 

The shock is advancing with almost three times the speed of sound!

Behind the shock the induced velocity is  $u_p = 756.2 \, m/s \Rightarrow$  supersonic flow  $(a_2 = 562.1 \, m/s)$ 

May be calculated by formulas 7.13, 7.16, 7.10, 7.11 or by using Table A.2 for stationary normal shock ( $u_1 = W$ ,  $u_2 = W - u_p$ )

**Note!**  $h_{O_1} \neq h_{O_2}$ 

constant total enthalpy is only valid for stationary shocks!

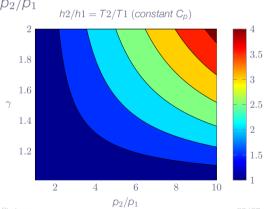
shock is uniquely defined by pressure ratio  $p_2/p_1$ 

$$u_{1} = 0$$

$$h_{o_{1}} = h_{1} + \frac{1}{2}u_{1}^{2} = h_{1}$$

$$h_{o_{2}} = h_{2} + \frac{1}{2}u_{2}^{2}$$

$$h_{2} > h_{1} \Rightarrow h_{o_{2}} > h_{o_{1}}$$



Gas/Vapor	Ratio of specific heats $(\gamma)$	Gas constant
Acatulana	1.23	319
Acetylene		
Air (standard)	1.40	287
Ammonia	1.31	530
Argon	1.67	208
Benzene	1.12	100
Butane	1.09	143
Carbon Dioxide	1.29	189
Carbon Disulphide	1.21	120
Carbon Monoxide	1.40	297
Chlorine	1.34	120
Ethane	1.19	276
Ethylene	1.24	296
Helium	1.67	2080
Hydrogen	1.41	4120
Hydrogen chloride	1.41	230
Methane	1,30	518
Natural Gas (Methane)	1.27	500
Nitric oxide	1.39	277
Nitrogen	1.40	297
Nitrous oxide	1,27	180
Oxygen	1.40	260
Propane	1.13	189
Steam (water)	1.32	462
Sulphur dioxide	1.29	130