Compressible Flow - TME085 Lecture 5

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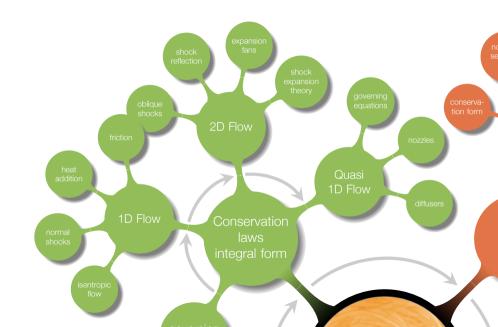
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Chapter 4 Oblique Shocks and Expansion Waves

Overview

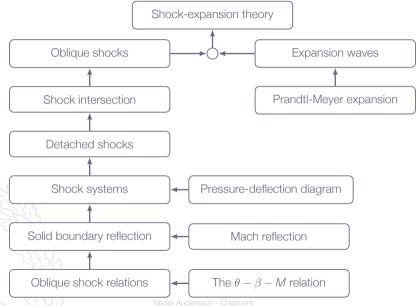


Learning Outcomes

- 4 Present at least two different formulations of the governing equations for compressible flows and explain what basic conservation principles they are based on
- 7 Explain why entropy is important for flow discontinuities
- 8 Derive (marked) and apply (all) of the presented mathematical formulae for classical gas dynamics
 - b normal shocks*
 - e oblique shocks in 2D*
 - f shock reflection at solid walls*
 - g contact discontinuities
 - h Prandtl-Meyer expansion fans in 2D
 - I detached blunt body shocks, nozzle flows
- 9 Solve engineering problems involving the above-mentioned phenomena (8a-8k)

why do we get normal shocks in some cases and oblique shocks in other?

Roadmap - Oblique Shocks and Expansion Waves



Motivation

Come on, two-dimensional flow, really?! Why not three-dimensional?

the normal shocks studied in chapter 3 are a special casees of the more general oblique shock waves that may be studied in two dimensions

in two dimensions, we can still analyze shock waves using a pen-and-paper approach

many practical problems or subsets of problems may be analyzed in two-dimensions

by going from one to two dimensions we will be able to introduce physical processes important for compressible flows

Oblique Shocks and Expansion Waves - Assumptions

- 1. Supersonic
- 2. Steady-state
- 3. Two-dimensional
- 4. Inviscid flow (no wall friction)

In real flow, viscosity is non-zero ⇒ boundary layers

For high-Reynolds-number flows, boundary layers are thin \Rightarrow inviscid theory still relevant!

Mach Wave

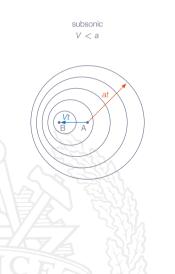
Sound waves emitted from A (speed of sound a)

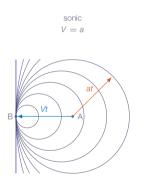


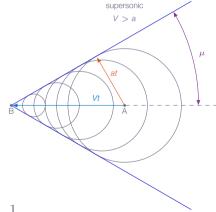


Mach Waves

A Mach wave is an infinitely weak oblique shock



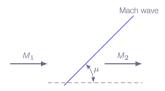




$$\sin \mu = \frac{\mathbf{a}t}{Vt} = \frac{\mathbf{a}}{V} = \frac{1}{M}$$

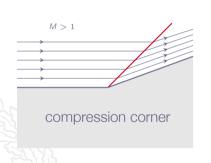
Mach Wave

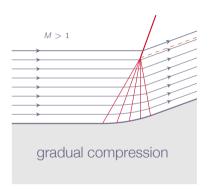
A Mach wave is an infinitely weak oblique shock

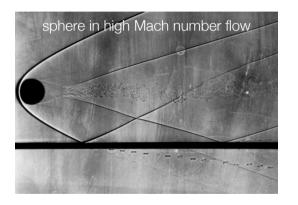


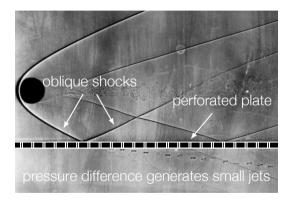
No substantial changes of flow properties over a single Mach wave $M_1>1.0$ and $M_1\approx M_2$ Isentropic

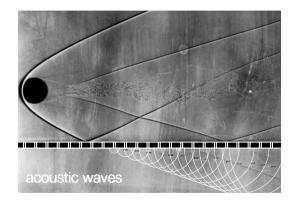
Oblique Shocks

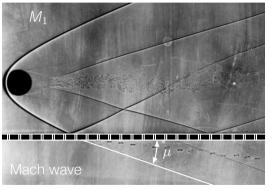






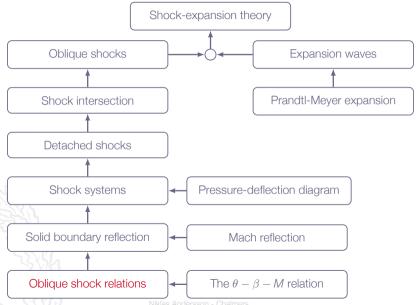






$$\mu = 19^{\circ} \Rightarrow M_1 = \frac{1}{\sin \mu} \approx 3.1$$

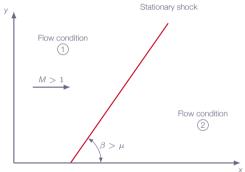
Roadmap - Oblique Shocks and Expansion Waves



Chapter 4.3 Oblique Shock Relations

Oblique Shocks

Two-dimensional steady-state flow

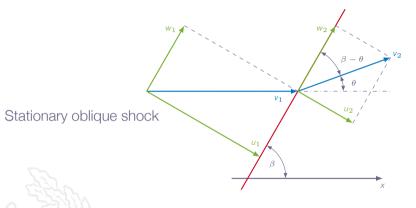


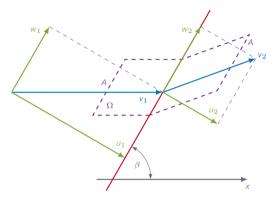
Significant changes of flow properties from 1 to 2

$$M_1>1.0,\,\beta>\mu,\,$$
 and $M_1\neq M_2$

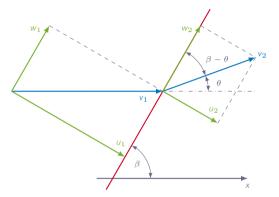
Not isentropic

Oblique Shocks





Two-dimensional steady-state flow
Control volume aligned with flow stream lines



Velocity notations:

$$M_{n_1} = \frac{u_1}{a_1} = M_1 \sin(\beta)$$

$$M_1 = \frac{v_2}{a_2}$$

$$M_{n_2} = \frac{u_2}{a_2} = M_2 \sin(\beta - \theta)$$

$$M_2 = \frac{v_2}{a_2}$$

Conservation of mass:

$$\frac{d}{dt} \iiint_{\Omega} \rho d\mathcal{V} + \iint_{\partial \Omega} \rho \mathbf{v} \cdot \mathbf{n} dS = 0$$

Mass conservation for control volume Ω :

$$0 - \rho_1 u_1 A + \rho_2 u_2 A = 0 \Rightarrow$$

$$\rho_1 u_1 = \rho_2 u_2$$

Conservation of momentum:

$$\frac{d}{dt} \iiint_{\Omega} \rho \mathbf{v} d\mathcal{V} + \iint_{\partial \Omega} \left[\rho(\mathbf{v} \cdot \mathbf{n}) \mathbf{v} + \rho \mathbf{n} \right] dS = \iiint_{\Omega} \rho \mathbf{f} d\mathcal{V}$$

Momentum in shock-normal direction:

$$0 - (\rho_1 u_1^2 + \rho_1)A + (\rho_2 u_2^2 + \rho_2)A = 0 \Rightarrow$$

$$\rho_1 u_1^2 + \rho_1 = \rho_2 u_2^2 + \rho_2$$

Momentum in shock-tangential direction:

$$0 - \rho_1 \mathbf{u}_1 \mathbf{w}_1 \mathbf{A} + \rho_2 \mathbf{u}_2 \mathbf{w}_2 \mathbf{A} = 0 \Rightarrow$$

$$\mathbf{w}_1 = \mathbf{w}_2$$

Conservation of energy:

$$\frac{d}{dt} \iiint_{\Omega} \rho \mathbf{e}_{o} d\mathcal{V} + \iint_{\partial\Omega} \left[\rho h_{o} \mathbf{v} \cdot \mathbf{n} \right] dS = \iiint_{\Omega} \rho \mathbf{f} \cdot \mathbf{v} d\mathcal{V}$$

Energy equation applied to the control volume Ω :

$$0 - \rho_1 u_1 [h_1 + \frac{1}{2} (u_1^2 + w_1^2)] A + \rho_2 u_2 [h_2 + \frac{1}{2} (u_2^2 + w_2^2)] A = 0 \Rightarrow$$

$$h_1 + \frac{1}{2}u_1^2 = h_2 + \frac{1}{2}u_2^2$$

We can use the same equations as for normal shocks if we replace M_1 with M_{n_1} and M_2 with M_{n_2}

$$M_{n_2}^2 = \frac{M_{n_1}^2 + [2/(\gamma - 1)]}{[2\gamma/(\gamma - 1)]M_{n_1}^2 - 1}$$

Ratios such as ρ_2/ρ_1 , ρ_2/ρ_1 , and T_2/T_1 can be calculated using the relations for normal shocks with M_1 replaced by M_{n_1}

What about ratios involving stagnation flow properties, can we use the ones previously derived for normal shocks?

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The shock process is adiabatic and thus total temperature is not effected by the shock $\Rightarrow T_{o_2} = T_{o_1}$



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What about the total pressure?

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The shock process is adiabatic and thus total temperature is not effected by the shock $\Rightarrow T_{o_2} = T_{o_1}$

What about the total pressure?

$$s_2 - s_1 = C_\rho \ln \left(\frac{T_{o_2}}{T_{o_1}} \right) - R \ln \left(\frac{p_{o_2}}{p_{o_1}} \right) = \{ T_{o_2} = T_{o_1} \} = -R \ln \left(\frac{p_{o_2}}{p_{o_1}} \right)$$

entropy is a thermodynamic flow property and s_2-s_1 is dictated by the shock strength and thus the total pressure ratio is a function of the shock-normal Mach number

Note! total pressure is always calculated using the flow Mach number, not the shock-normal Mach number

However, the ratio p_{o_2}/p_{o_1} may be calculated using the shock-normal Mach number

So, be careful when using relations derived for normal shocks for oblique shocks when it comes to total flow conditions...

$$p_{o_2}/p_{o_1}$$
 is calculated as: $\frac{p_{o_2}}{p_{o_1}} = \frac{p_{o_2}}{p_2} \frac{p_2}{p_1} \frac{p_1}{p_{o_1}}$

where

1.
$$\frac{\rho_{o_2}}{\rho_2} = f(M_2), \frac{\rho_2}{\rho_1} = f(M_{n_1}), \text{ and } \frac{\rho_1}{\rho_{o_1}} = f(M_1)$$

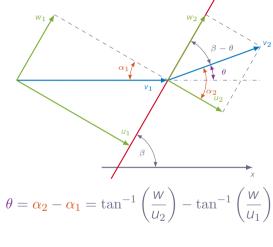
or alternatively

2.
$$\frac{\rho_{o_2}}{\rho_2} = f(M_{n_2}), \frac{\rho_2}{\rho_1} = f(M_{n_1}), \text{ and } \frac{\rho_1}{\rho_{o_1}} = f(M_{n_1})$$

Note! in the second case the total pressures are **not** the true total pressures of the flow and therefore it is suggested to use the first approach

Deflection Angle (for the interested)





$$\frac{\partial \theta}{\partial W} = \frac{u_2}{W^2 + u_2^2} - \frac{u_1}{W^2 + u_1^2}$$

Deflection Angle (for the interested)



$$\frac{\partial \theta}{\partial w} = \frac{u_2}{w^2 + u_2^2} - \frac{u_1}{w^2 + u_1^2} = 0 \Rightarrow$$

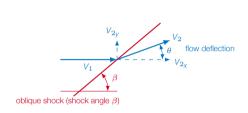
$$\frac{u_2(w^2 + u_1^2) - u_1(w^2 + u_2^2)}{(w^2 + u_2^2)(w^2 + u_1^2)} = 0 \Rightarrow \frac{(u_2 - u_1)(w^2 - u_1u_2)}{(w^2 + u_2^2)(w^2 + u_1^2)} = 0$$

Two solutions:

- $ightharpoonup u_2 = u_1$ (no deflection)
- $ightharpoonup w^2 = u_1 u_2$ (max deflection)

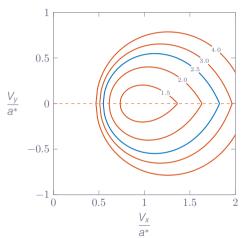
Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number



No deflection cases:

- normal shock(reduced shock-normal velocity)
- Mach wave (unchanged shock-normal velocity)

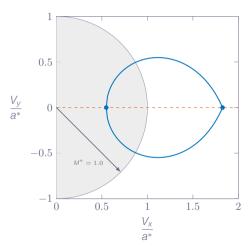


Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

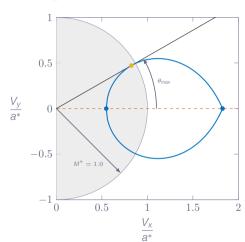
$$M^* = \frac{\sqrt{V_X^2 + V_y^2}}{a^*}$$

Solutions to the left of the sonic line are subsonic



Graphical representation of all possible deflection angles for a specific Mach number

It is not possible to deflect the flow more than θ_{max}

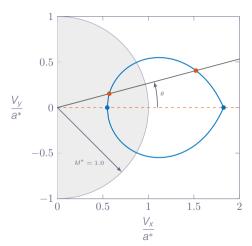


Graphical representation of all possible deflection angles for a specific Mach number

For each deflection angle $\theta < \theta_{max}$, there are two solutions

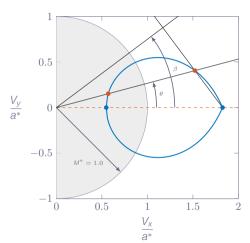
- strong shock solution
- weak shock solution

Weak shocks give lower losses and therefore the preferred solution



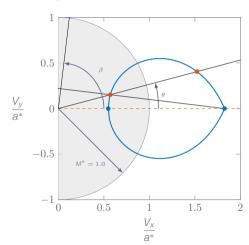
Graphical representation of all possible deflection angles for a specific Mach number

The shock polar can be used to calculate the shock angle β for a given deflection angle θ

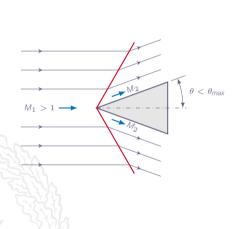


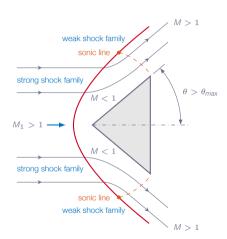
Graphical representation of all possible deflection angles for a specific Mach number

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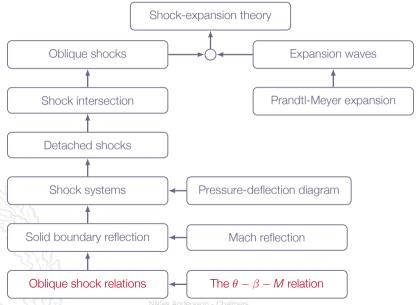


Flow Deflection





Roadmap - Oblique Shocks and Expansion Waves



The θ - β -M Relation

It can be shown that

$$\tan \theta = 2 \cot \beta \left(\frac{M_1^2 \sin^2 \beta - 1}{M_1^2 (\gamma + \cos 2\beta) + 2} \right)$$

which is the θ - β -M relation

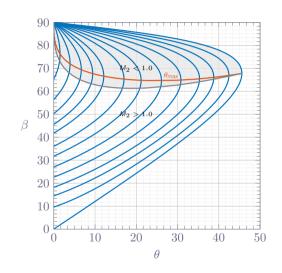
The θ - β -Mach Relation

A relation between:

- ightharpoonup flow deflection angle θ
- ightharpoonup shock angle β
- ▶ upstream flow Mach number M₁

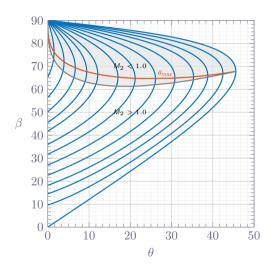
$$\tan(\theta) = 2\cot(\beta) \left(\frac{M_1^2 \sin^2(\beta) - 1}{M_1^2 (\gamma + \cos(2\beta)) + 2} \right)$$

Note! in general there are two solutions for a given M_1 (or none)



The θ - β -Mach Relation

- There is a small region where we may find weak shock solutions for which $M_2 < 1$
- In most cases weak shock solutions have $M_2 > 1$
- ▶ Strong shock solutions always have $M_2 < 1$
- In practical situations, weak shock solutions are most common
- Strong shock solution may appear in special situations due to high back pressure, which forces $M_2 < 1$



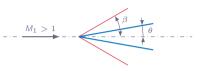
The θ - β -M Relation

Note! In Chapter 3 we learned that the Mach number always reduces to subsonic values behind a shock. This is true for normal shocks (shocks that are normal to the flow direction). It is also true for oblique shocks if looking in the shock-normal direction.

The θ - β -M Relation - Wedge Flow

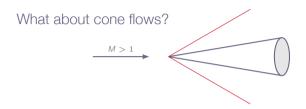
Wedge flow oblique shock analysis:

- 1. θ - β -M relation $\Rightarrow \beta$ for given M_1 and θ
- 2. β gives M_{n_1} according to: $M_{n_1} = M_1 \sin(\beta)$
- 3. normal shock formula with M_{n_1} instead of $M_1 \Rightarrow M_{n_2}$ (instead of M_2)
- 4. M_2 given by $M_2 = M_{n_2}/\sin(\beta \theta)$
- 5. normal shock formula with M_{n_1} instead of $M_1 \Rightarrow \rho_2/\rho_1, \rho_2/\rho_1$, etc
- 6. upstream conditions + ρ_2/ρ_1 , ρ_2/ρ_1 , etc \Rightarrow downstream conditions



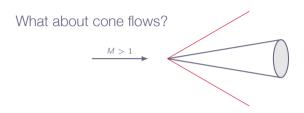
Chapter 4.4 Supersonic Flow over Wedges and Cones

Supersonic Flow over Wedges and Cones



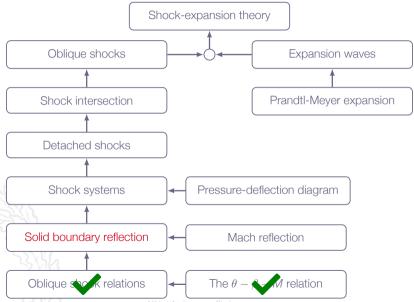
- Similar to wedge flow, we do get a constant-strength shock wave, attached at the cone tip (or else a detached curved shock)
- The attached shock is also cone-shaped

Supersonic Flow over Wedges and Cones



- ► The flow condition immediately downstream of the shock is uniform
- ► However, downstream of the shock the streamlines are curved and the flow varies in a more complex manner (3D relieving effect as R increases there is more and more space around cone for the flow)
- \triangleright β for cone shock is always smaller than that for wedge shock, if M_1 is the same

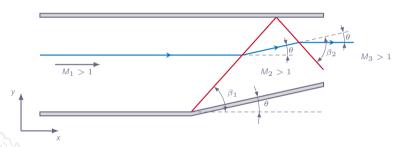
Roadmap - Oblique Shocks and Expansion Waves



Chapter 4.6
Regular Reflection from a Solid
Boundary

Shock Reflection

Regular reflection of oblique shock at solid wall (see example 4.10)



Assumptions:

- steady-state inviscid flow
- weak shocks

Shock Reflection

first shock:

upstream condition:

 $M_1 > 1$, flow in x-direction

downstream condition:

weak shock $\Rightarrow M_2 > 1$ deflection angle θ shock angle β_1

second shock:

upstream condition:

same as downstream condition of first shock

downstream condition:

weak shock $\Rightarrow M_3 > 1$ deflection angle θ shock angle β_2

Shock Reflection

Solution:

first shock:

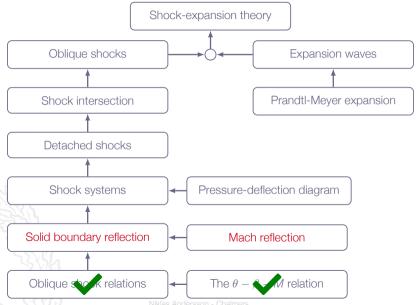
- \triangleright β_1 calculated from θ - β -M relation for specified θ and M_1 (weak solution)
- ▶ flow condition 2 according to formulas for normal shocks $(M_{n_1} = M_1 \sin(\beta_1))$ and $M_{n_2} = M_2 \sin(\beta_1 \theta)$

second shock:

- $\nearrow \geqslant \beta_2$ calculated from θ - β -M relation for specified θ and M_2 (weak solution)
- flow condition 3 according to formulas for normal shocks $(M_{n_2} = M_2 \sin(\beta_2))$ and $M_{n_3} = M_3 \sin(\beta_2 \theta)$

 \Rightarrow complete description of flow and shock waves (angle between upper wall and second shock: $\Phi=\beta_2-\theta$)

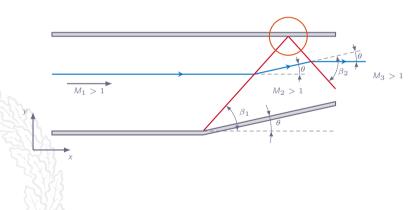
Roadmap - Oblique Shocks and Expansion Waves



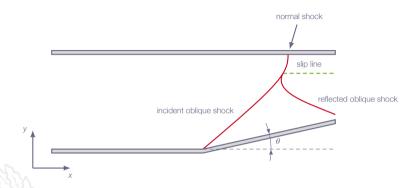
Chapter 4.11 Mach Reflection

Regular Shock Reflection

Regular reflection possible if both primary and reflected shocks are weak (see θ - β -M relation)

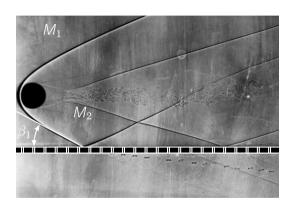


Mach Reflection



Mach reflection:

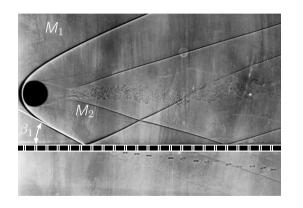
- appears when regular reflection is not possible
- more complex flow than for a regular reflection
- no analytic solution numerical solution necessary



$$\theta_1 = f(M_1, \beta_1), \quad M_2 = f(M_1, \theta_1, \beta_1)$$

 $M_1 > M_2$

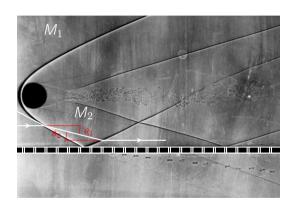
 $M_2 > 1.0$



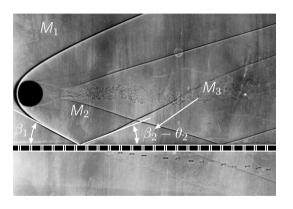
$$\beta_1 = 28^{\circ}$$

$$M_1 = 3.1$$

$$\Rightarrow \theta_1 \approx 11.2^{\circ}, \quad M_2 \approx 2.5$$







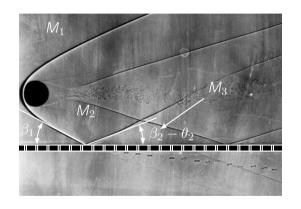
$$M_1 > M_2 > M_3$$

$$M_3 > 1.0$$

$$\beta_2 > \beta_1$$

$$\beta_2 = f(M_2, \theta_2), \quad M_3 = f(M_2, \theta_2, \beta_2)$$

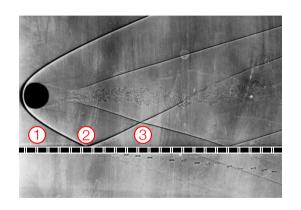
Note! Shock wave reflection at solid wall is not specular



$$\theta_2 = 11.2^{\circ}$$

$$M_2 = 2.5$$

$$\Rightarrow \beta_2 \approx 33^{\circ}, \quad M_3 \approx 2.0$$



$$\frac{\rho_3}{\rho_1} = \frac{\rho_2}{\rho_1} \frac{\rho_3}{\rho_2} \approx 4.52$$

$$\frac{T_3}{T_1} = \frac{T_2}{T_1} \frac{T_3}{T_2} \approx 1.57$$