

# Compressible Flow - TME085

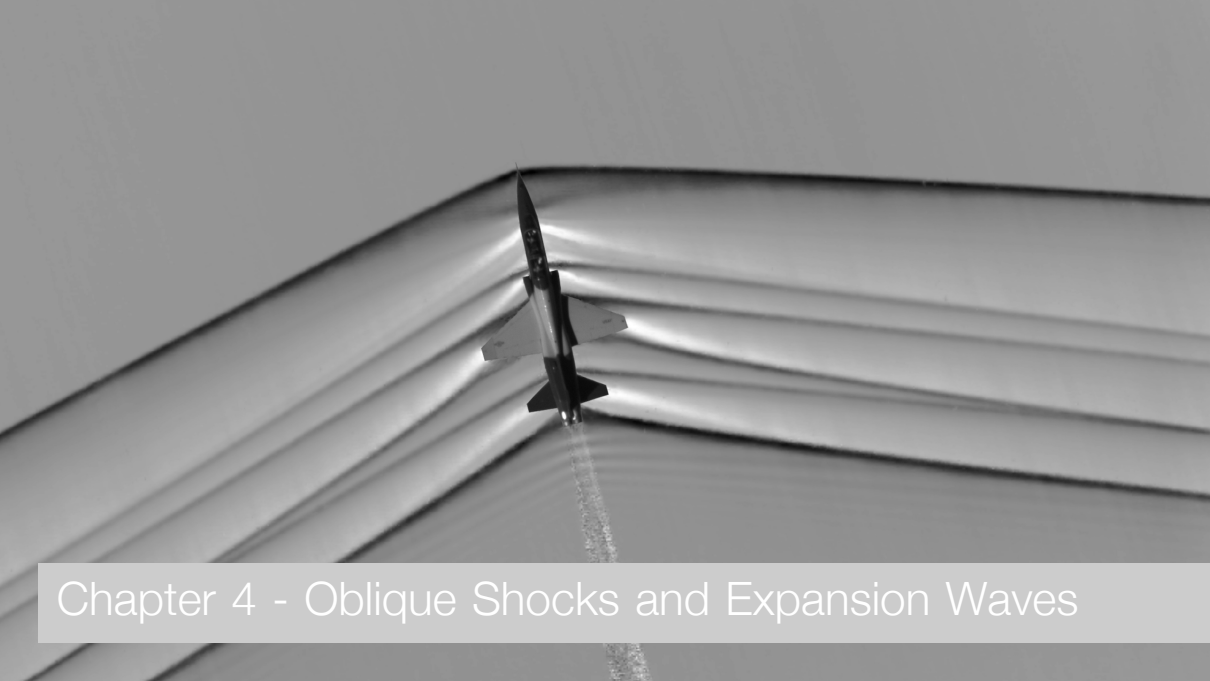
## Lecture 5

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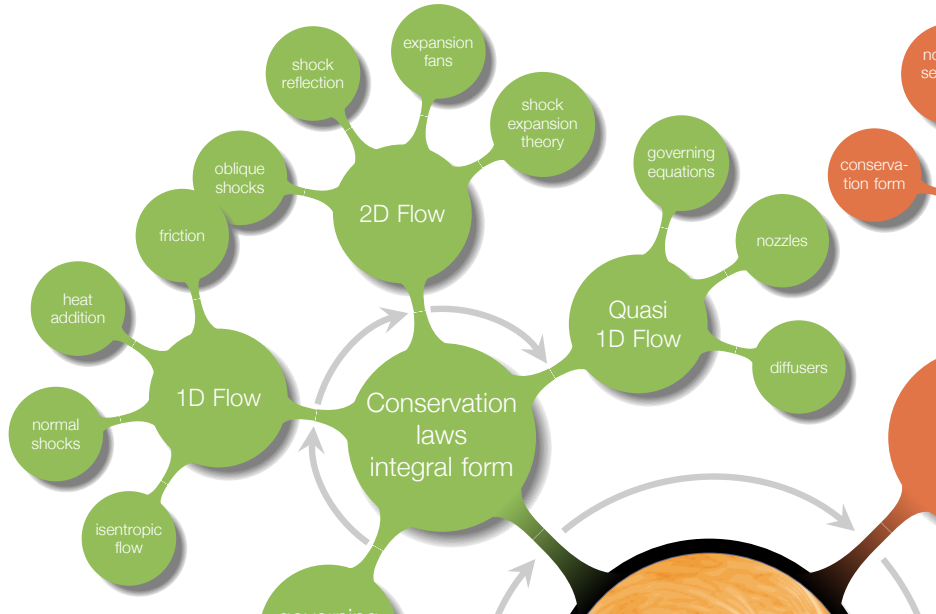
`niklas.andersson@chalmers.se`





## Chapter 4 - Oblique Shocks and Expansion Waves

# Overview

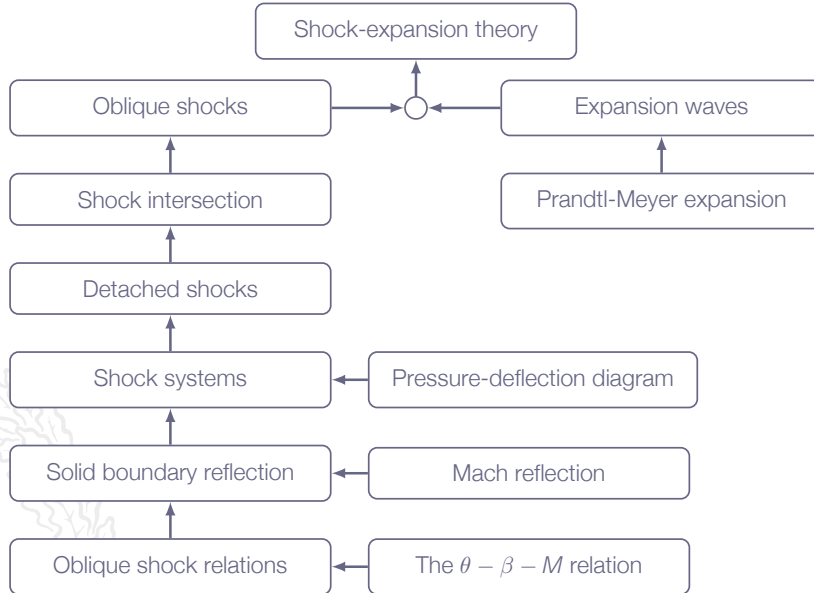


# Learning Outcomes

- 4 **Present** at least two different formulations of the governing equations for compressible flows and **explain** what basic conservation principles they are based on
- 7 **Explain** why entropy is important for flow discontinuities
- 8 **Derive** (marked) and **apply** (all) of the presented mathematical formulae for classical gas dynamics
  - b normal shocks\*
  - e oblique shocks in 2D\*
  - f shock reflection at solid walls\*
  - g contact discontinuities
  - h Prandtl-Meyer expansion fans in 2D
  - i detached blunt body shocks, nozzle flows
- 9 **Solve** engineering problems involving the above-mentioned phenomena (8a-8k)

*why do we get normal shocks in some cases and oblique shocks in other?*

# Roadmap - Oblique Shocks and Expansion Waves



# Motivation

Come on, two-dimensional flow, really?! Why not three-dimensional?

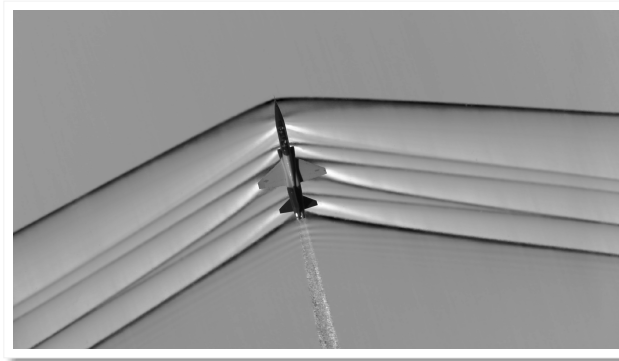
the normal shocks studied in chapter 3 are a special case of the more general oblique shock waves that may be studied in two dimensions

in two dimensions, we can still analyze shock waves using a pen-and-paper approach

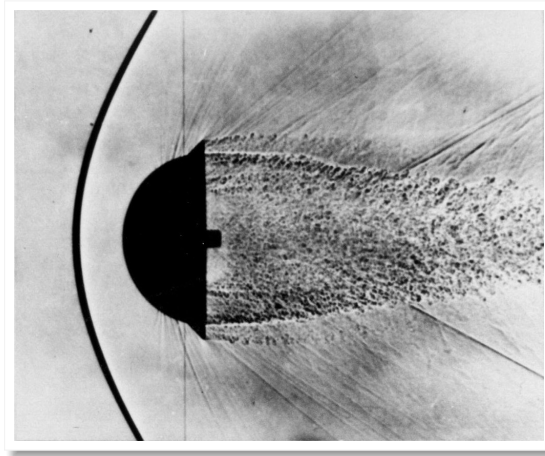
many practical problems or subsets of problems may be analyzed in two-dimensions

by going from one to two dimensions we will be able to introduce physical processes important for compressible flows

# Oblique Shocks and Expansion Waves



# Oblique Shocks and Expansion Waves





# Oblique Shocks and Expansion Waves - Assumptions

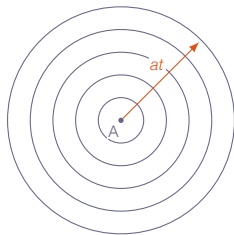
1. Supersonic
2. Steady-state
3. Two-dimensional
4. Inviscid flow (no wall friction)

In real flow, viscosity is non-zero  $\Rightarrow$  boundary layers

For high-Reynolds-number flows, boundary layers are thin  $\Rightarrow$  inviscid theory still relevant!

# Mach Wave

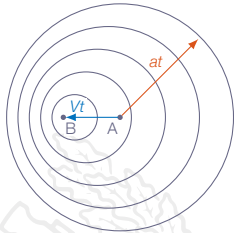
Sound waves emitted from A (speed of sound  $a$ )



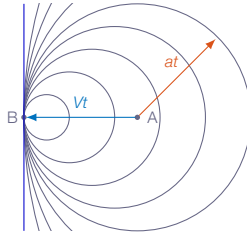
# Mach Waves

A Mach wave is an infinitely weak oblique shock

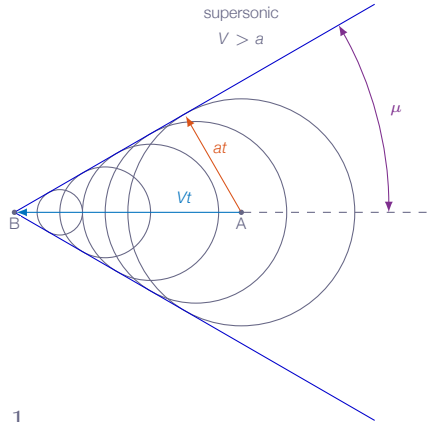
subsonic  
 $V < a$



sonic  
 $V = a$



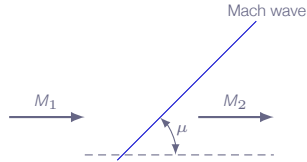
supersonic  
 $V > a$



$$\sin \mu = \frac{at}{Vt} = \frac{a}{V} = \frac{1}{M}$$

# Mach Wave

A Mach wave is an infinitely weak oblique shock

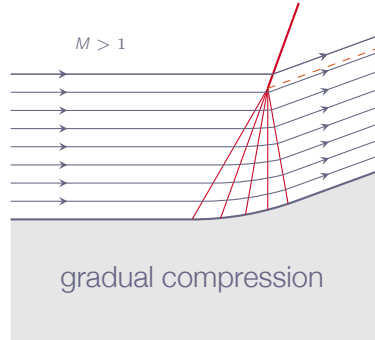
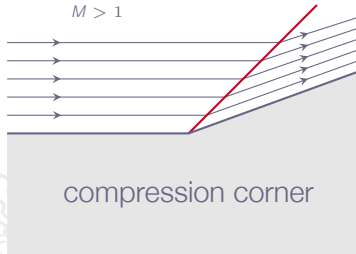


No substantial changes of flow properties over a single Mach wave

$M_1 > 1.0$  and  $M_1 \approx M_2$

Isentropic

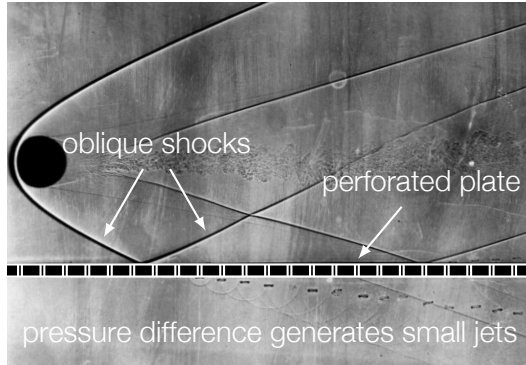
# Oblique Shocks



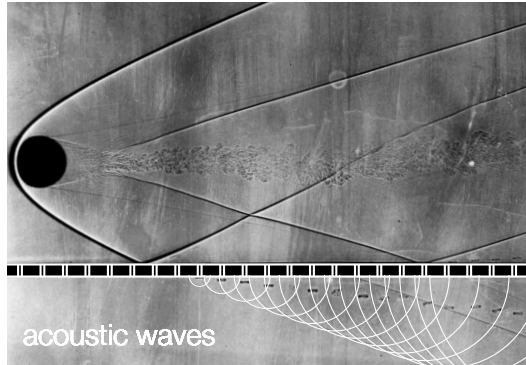
# Oblique Shocks and Mach Waves



# Oblique Shocks and Mach Waves

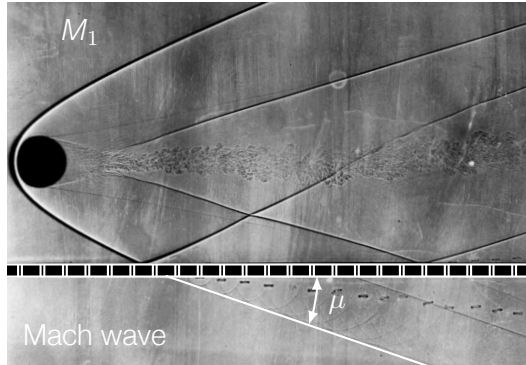


# Oblique Shocks and Mach Waves



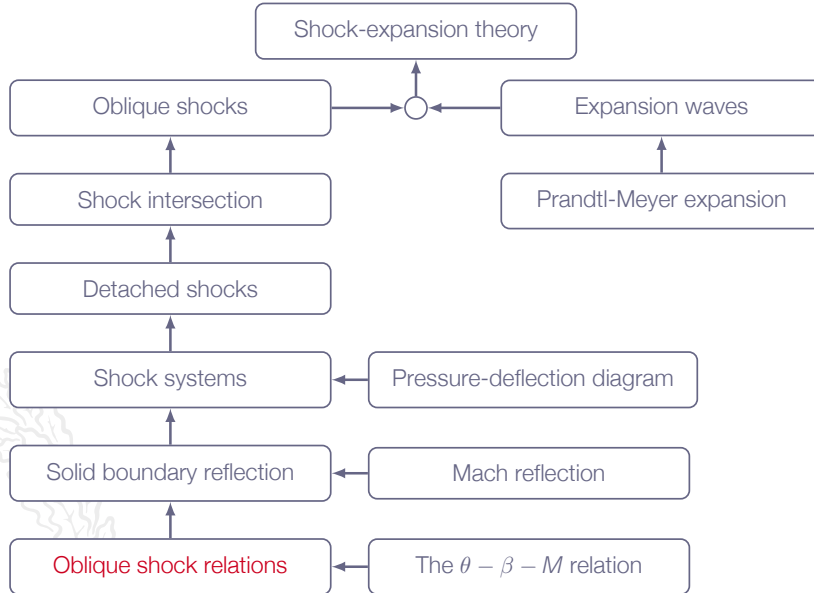


# Oblique Shocks and Mach Waves



$$\mu = 19^\circ \Rightarrow M_1 = \frac{1}{\sin \mu} \approx 3.1$$

# Roadmap - Oblique Shocks and Expansion Waves



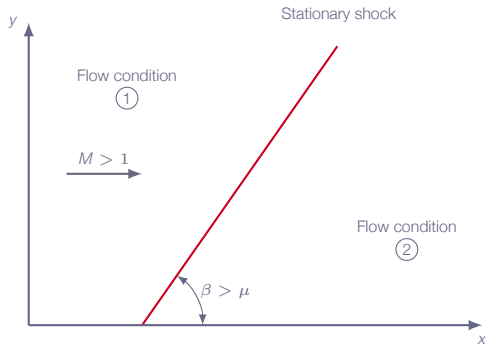
# Chapter 4.3

## Oblique Shock Relations



# Oblique Shocks

Two-dimensional steady-state flow



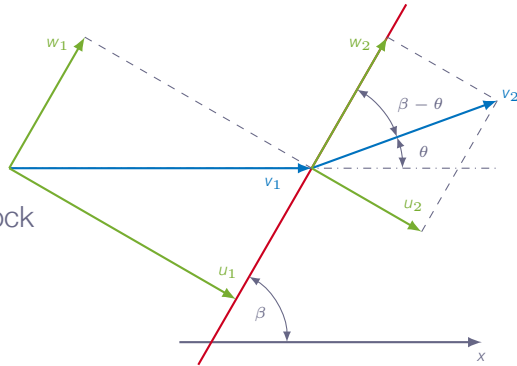
Significant changes of flow properties from 1 to 2

$M_1 > 1.0$ ,  $\beta > \mu$ , and  $M_1 \neq M_2$

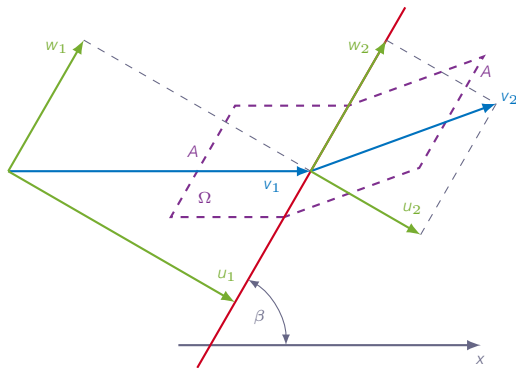
Not isentropic

# Oblique Shocks

Stationary oblique shock



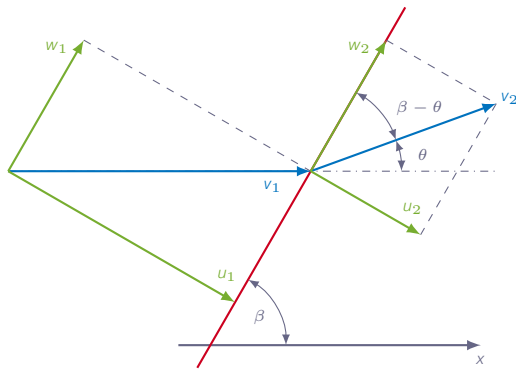
# Oblique Shock Relations



Two-dimensional steady-state flow

Control volume aligned with flow stream lines

# Oblique Shock Relations



Velocity notations:

$$M_{n1} = \frac{u_1}{a_1} = M_1 \sin(\beta)$$

$$M_1 = \frac{v_1}{a_1}$$

$$M_{n2} = \frac{u_2}{a_2} = M_2 \sin(\beta - \theta)$$

$$M_2 = \frac{v_2}{a_2}$$

# Oblique Shock Relations

Conservation of mass:

$$\frac{d}{dt} \iiint_{\Omega} \rho d\mathcal{V} + \oiint_{\partial\Omega} \rho \mathbf{v} \cdot \mathbf{n} dS = 0$$

Mass conservation for control volume  $\Omega$ :

$$0 - \rho_1 u_1 A + \rho_2 u_2 A = 0 \Rightarrow$$

$$\rho_1 u_1 = \rho_2 u_2$$



# Oblique Shock Relations

Conservation of momentum:

$$\frac{d}{dt} \iiint_{\Omega} \rho \mathbf{v} d\mathcal{V} + \oint_{\partial\Omega} [\rho(\mathbf{v} \cdot \mathbf{n})\mathbf{v} + p\mathbf{n}] dS = \iiint_{\Omega} \rho \mathbf{f} d\mathcal{V}$$

Momentum in shock-normal direction:

$$0 - (\rho_1 u_1^2 + p_1)A + (\rho_2 u_2^2 + p_2)A = 0 \Rightarrow$$

$$\rho_1 u_1^2 + p_1 = \rho_2 u_2^2 + p_2$$

# Oblique Shock Relations

Momentum in shock-tangential direction:

$$0 - \rho_1 u_1 w_1 A + \rho_2 u_2 w_2 A = 0 \Rightarrow$$

$$w_1 = w_2$$



# Oblique Shock Relations

Conservation of energy:

$$\frac{d}{dt} \iiint_{\Omega} \rho e_o d\mathcal{V} + \iint_{\partial\Omega} [\rho h_o \mathbf{v} \cdot \mathbf{n}] dS = \iiint_{\Omega} \rho \mathbf{f} \cdot \mathbf{v} d\mathcal{V}$$

Energy equation applied to the control volume  $\Omega$ :

$$0 - \rho_1 u_1 \left[ h_1 + \frac{1}{2} (u_1^2 + w_1^2) \right] A + \rho_2 u_2 \left[ h_2 + \frac{1}{2} (u_2^2 + w_2^2) \right] A = 0 \Rightarrow$$

$$h_1 + \frac{1}{2} u_1^2 = h_2 + \frac{1}{2} u_2^2$$

# Oblique Shock Relations

We can use the same equations as for normal shocks if we replace  $M_1$  with  $M_{n_1}$  and  $M_2$  with  $M_{n_2}$

$$M_{n_2}^2 = \frac{M_{n_1}^2 + [2/(\gamma - 1)]}{[2\gamma/(\gamma - 1)]M_{n_1}^2 - 1}$$

Ratios such as  $\rho_2/\rho_1$ ,  $p_2/p_1$ , and  $T_2/T_1$  can be calculated using the relations for normal shocks with  $M_1$  replaced by  $M_{n_1}$

# Oblique Shock Relations

What about ratios involving stagnation flow properties, can we use the ones previously derived for normal shocks?



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The shock process is adiabatic and thus total temperature is not effected by the shock  $\Rightarrow T_{o_2} = T_{o_1}$



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What about ratios involving stagnation flow properties, can we use the ones previously derived for normal shocks?

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What about the total pressure?



# Oblique Shock Relations

What about ratios involving stagnation flow properties, can we use the ones previously derived for normal shocks?

The shock process is adiabatic and thus total temperature is not effected by the shock  $\Rightarrow T_{o2} = T_{o1}$

What about the total pressure?

$$s_2 - s_1 = C_p \ln \left( \frac{T_{o2}}{T_{o1}} \right) - R \ln \left( \frac{p_{o2}}{p_{o1}} \right) = \{T_{o2} = T_{o1}\} = -R \ln \left( \frac{p_{o2}}{p_{o1}} \right)$$

entropy is a thermodynamic flow property and  $s_2 - s_1$  is dictated by the shock strength and thus the total pressure ratio is a function of the shock-normal Mach number



# Oblique Shock Relations

**Note!** total pressure is always calculated using the flow Mach number, not the shock-normal Mach number

However, the ratio  $p_{o2}/p_{o1}$  may be calculated using the shock-normal Mach number

So, be careful when using relations derived for normal shocks for oblique shocks when it comes to total flow conditions...

# Oblique Shock Relations

$p_{o2}/p_{o1}$  is calculated as:  $\frac{p_{o2}}{p_{o1}} = \frac{p_{o2}}{p_2} \frac{p_2}{p_1} \frac{p_1}{p_{o1}}$

where

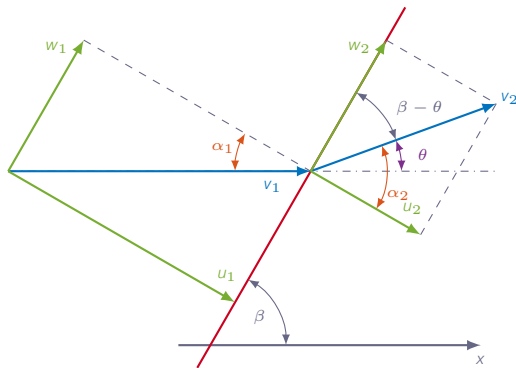
1.  $\frac{p_{o2}}{p_2} = f(M_2)$ ,  $\frac{p_2}{p_1} = f(M_{n1})$ , and  $\frac{p_1}{p_{o1}} = f(M_1)$

or alternatively

2.  $\frac{p_{o2}}{p_2} = f(M_{n2})$ ,  $\frac{p_2}{p_1} = f(M_{n1})$ , and  $\frac{p_1}{p_{o1}} = f(M_{n1})$

**Note!** in the second case the total pressures are **not** the true total pressures of the flow and therefore it is suggested to use the first approach

# Deflection Angle (for the interested)



$$\theta = \alpha_2 - \alpha_1 = \tan^{-1} \left( \frac{w}{u_2} \right) - \tan^{-1} \left( \frac{w}{u_1} \right)$$

$$\frac{\partial \theta}{\partial w} = \frac{u_2}{w^2 + u_2^2} - \frac{u_1}{w^2 + u_1^2}$$

## Deflection Angle (for the interested)



$$\frac{\partial \theta}{\partial w} = \frac{u_2}{w^2 + u_2^2} - \frac{u_1}{w^2 + u_1^2} = 0 \Rightarrow$$

$$\frac{u_2(w^2 + u_1^2) - u_1(w^2 + u_2^2)}{(w^2 + u_2^2)(w^2 + u_1^2)} = 0 \Rightarrow \frac{(u_2 - u_1)(w^2 - u_1 u_2)}{(w^2 + u_2^2)(w^2 + u_1^2)} = 0$$

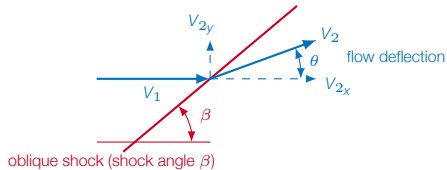
Two solutions:

$$u_2 = u_1 \text{ (no deflection)}$$

$$w^2 = u_1 u_2 \text{ (max deflection)}$$

# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

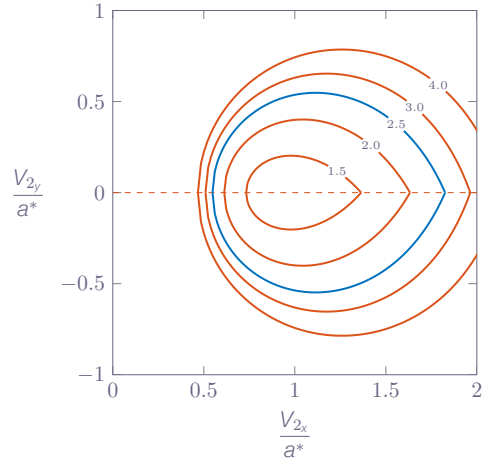


## Note!

In the shock polar,  $V_{2x}$  and  $V_{2y}$  are normalized by  $a^*$

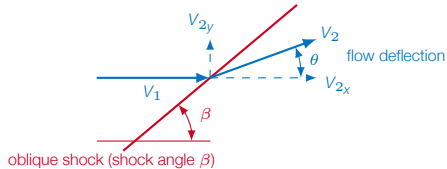
$a^*$  is a constant in a adiabatic flow

$a^*$  is not affected by shocks



# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number



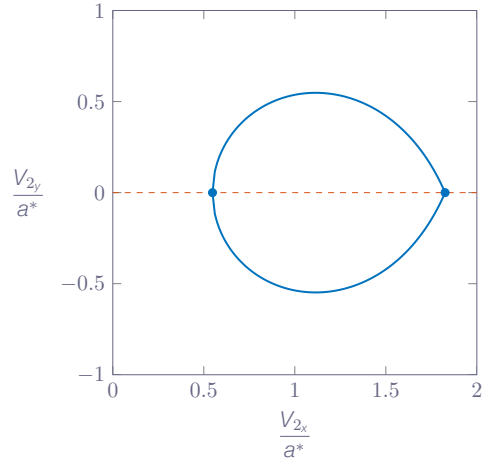
No deflection cases:

**normal shock**

(reduced shock-normal velocity)

**Mach wave**

(unchanged shock-normal velocity)



# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

$$M^* = \frac{\sqrt{V_{2x}^2 + V_{2y}^2}}{a^*}$$

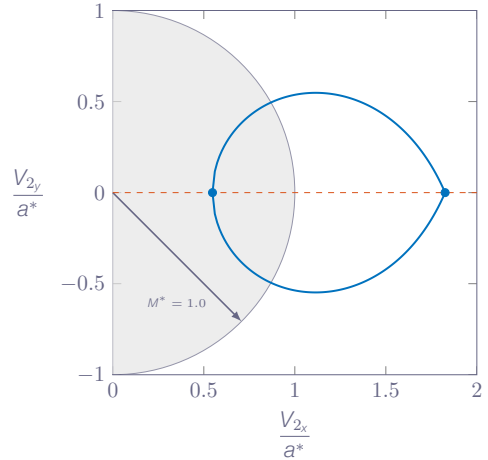
Solutions to the left of the sonic line are subsonic

## Recall

$$M^* = 1 \Leftrightarrow M = 1$$

$$M^* < 1 \Leftrightarrow M < 1$$

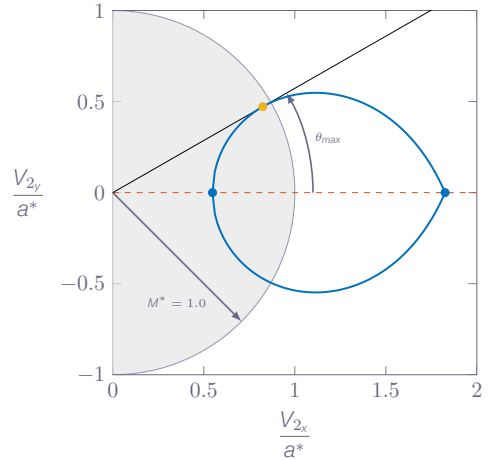
$$M^* > 1 \Leftrightarrow M > 1$$



# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

It is not possible to deflect the flow more than  $\theta_{max}$





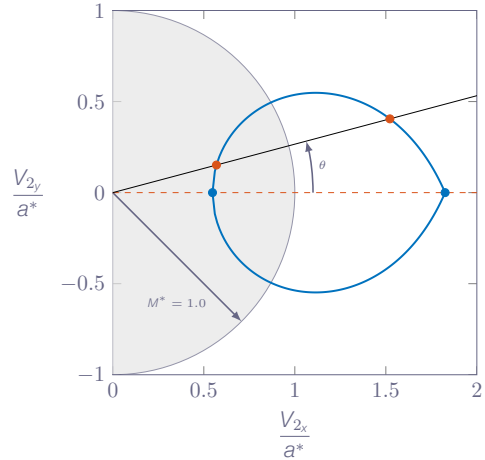
# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

For each deflection angle  $\theta < \theta_{max}$ , there are two solutions

1. **strong shock** solution
2. **weak shock** solution

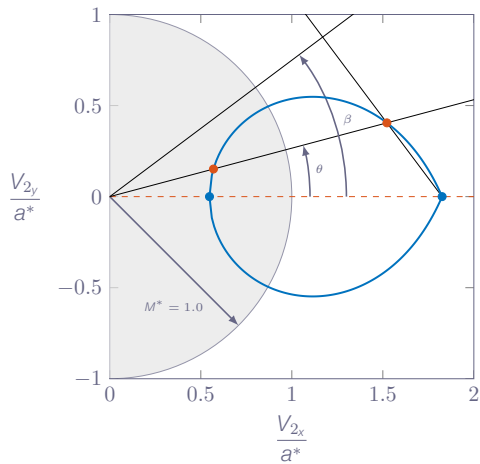
Weak shocks give lower losses and therefore the preferred solution



# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

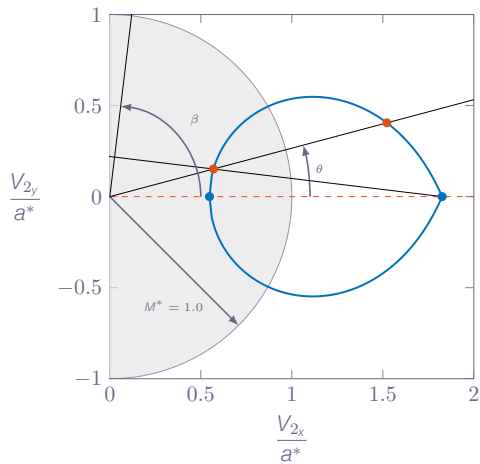
The shock polar can be used to calculate the shock angle  $\beta$  for a given deflection angle  $\theta$



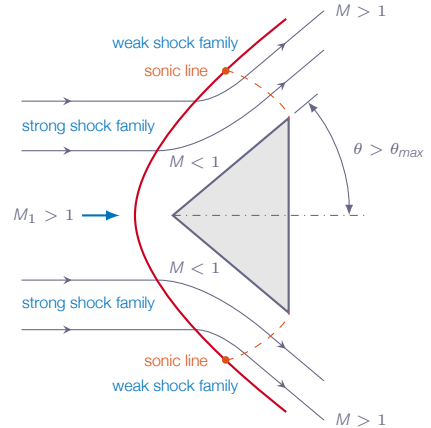
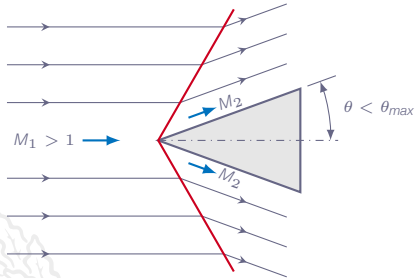
# Shock Polar

Graphical representation of all possible deflection angles for a specific Mach number

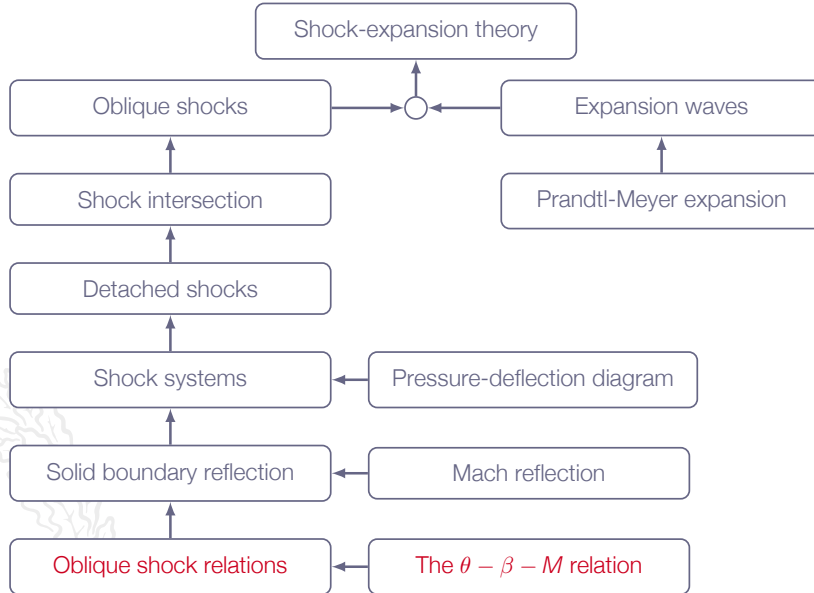
The shock polar can be used to calculate the shock angle  $\beta$  for a given deflection angle  $\theta$



# Flow Deflection



# Roadmap - Oblique Shocks and Expansion Waves

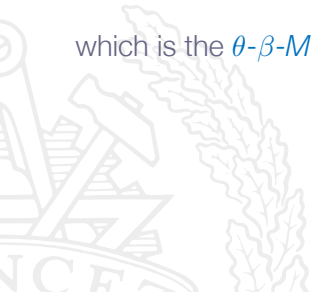


# The $\theta$ - $\beta$ - $M$ Relation

It can be shown that

$$\tan \theta = 2 \cot \beta \left( \frac{M_1^2 \sin^2 \beta - 1}{M_1^2 (\gamma + \cos 2\beta) + 2} \right)$$

which is the  $\theta$ - $\beta$ - $M$  relation



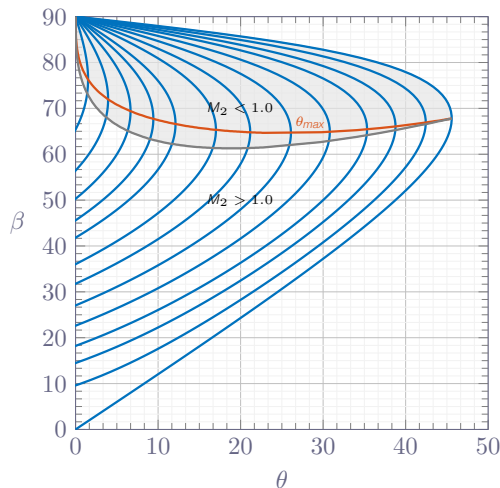
# The $\theta$ - $\beta$ -Mach Relation

A relation between:

1. flow deflection angle  $\theta$
2. shock angle  $\beta$
3. upstream flow Mach number  $M_1$

$$\tan(\theta) = 2 \cot(\beta) \left( \frac{M_1^2 \sin^2(\beta) - 1}{M_1^2(\gamma + \cos(2\beta)) + 2} \right)$$

**Note!** in general there are two solutions for a given  $M_1$  (or none)



# The $\theta$ - $\beta$ -Mach Relation

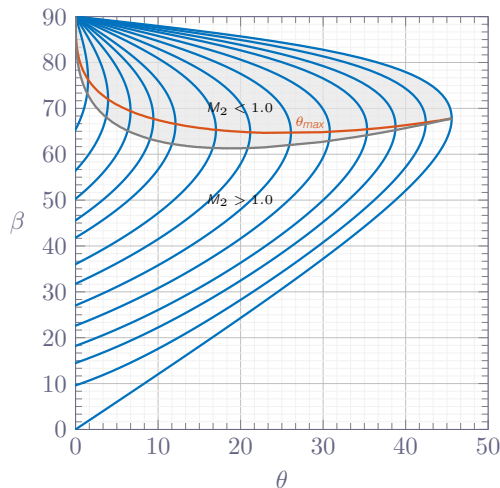
There is a small region where we may find weak shock solutions for which  $M_2 < 1$

In most cases weak shock solutions have  $M_2 > 1$

Strong shock solutions always have  $M_2 < 1$

In practical situations, **weak shock solutions are most common**

Strong shock solution may appear in special situations due to high back pressure, which forces  $M_2 < 1$





# The $\theta$ - $\beta$ - $M$ Relation

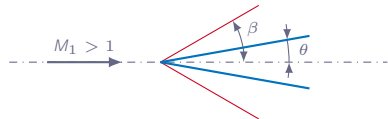
**Note!** In Chapter 3 we learned that the Mach number always reduces to subsonic values behind a shock. This is true for normal shocks (shocks that are normal to the flow direction). It is also true for oblique shocks if looking in the shock-normal direction.



# The $\theta$ - $\beta$ - $M$ Relation - Wedge Flow

Wedge flow oblique shock analysis:

1.  $\theta$ - $\beta$ - $M$  relation  $\Rightarrow \beta$  for given  $M_1$  and  $\theta$
2.  $\beta$  gives  $M_{n_1}$  according to:  $M_{n_1} = M_1 \sin(\beta)$
3. normal shock formula with  $M_{n_1}$  instead of  $M_1 \Rightarrow M_{n_2}$  (instead of  $M_2$ )
4.  $M_2$  given by  $M_2 = M_{n_2} / \sin(\beta - \theta)$
5. normal shock formula with  $M_{n_1}$  instead of  $M_1 \Rightarrow \rho_2/\rho_1, p_2/p_1$ , etc
6. upstream conditions +  $\rho_2/\rho_1, p_2/p_1$ , etc  $\Rightarrow$  downstream conditions





# Chapter 4.4

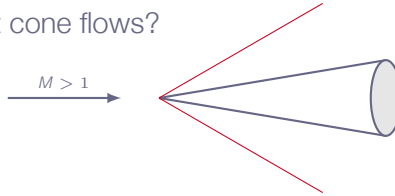
## Supersonic Flow over Wedges and Cones



# Supersonic Flow over Wedges and Cones



What about cone flows?



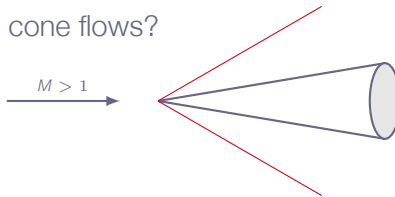
Similar to wedge flow, we do get a constant-strength shock wave, attached at the cone tip (or else a detached curved shock)

The attached shock is also cone-shaped

# Supersonic Flow over Wedges and Cones



What about cone flows?

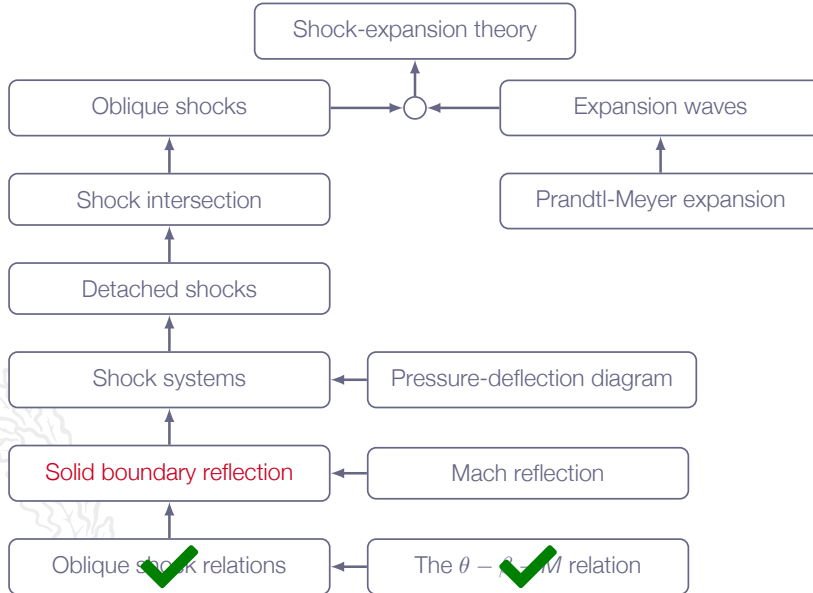


The flow condition immediately downstream of the shock is uniform

However, downstream of the shock the streamlines are curved and the flow varies in a more complex manner (3D relieving effect - as  $R$  increases there is more and more space around cone for the flow)

$\beta$  for cone shock is always smaller than that for wedge shock, if  $M_1$  is the same

# Roadmap - Oblique Shocks and Expansion Waves



# Chapter 4.6

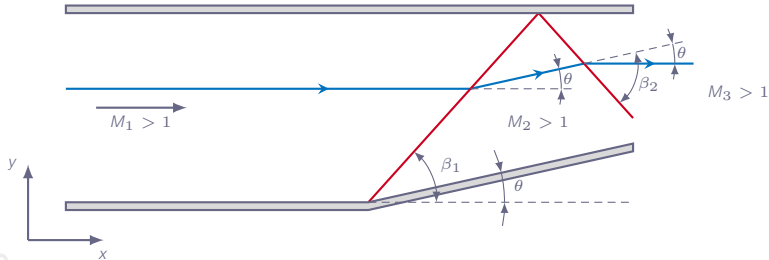
## Regular Reflection from a Solid Boundary



# Shock Reflection

## Regular reflection of oblique shock at solid wall

(see example 4.10)



Assumptions:

steady-state inviscid flow

weak shocks



# Shock Reflection

first shock

**upstream condition**

$$M_1 > 1$$

flow in x-direction

**downstream condition**

weak shock  $\Rightarrow M_2 > 1$

deflection angle  $\theta$

shock angle  $\beta_1$

second shock

**upstream condition**

downstream of first shock

**downstream condition**

weak shock  $\Rightarrow M_3 > 1$

deflection angle  $\theta$

shock angle  $\beta_2$

# Shock Reflection

Solution:

first shock:

1.  $\beta_1$  calculated from  $\theta$ - $\beta$ - $M$  relation for specified  $\theta$  and  $M_1$  (*weak solution*)
2. flow condition 2 according to formulas for normal shocks ( $M_{n_1} = M_1 \sin(\beta_1)$  and  $M_{n_2} = M_2 \sin(\beta_1 - \theta)$ )

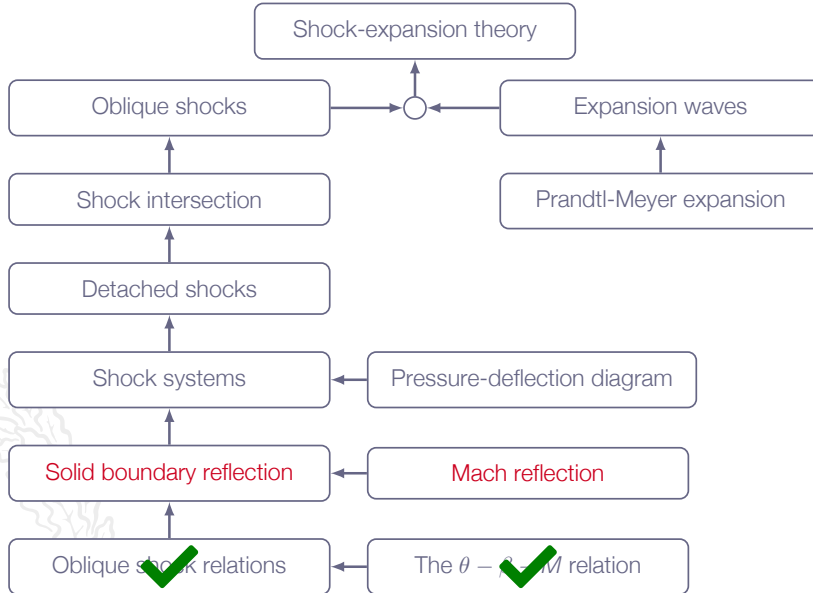
second shock:

1.  $\beta_2$  calculated from  $\theta$ - $\beta$ - $M$  relation for specified  $\theta$  and  $M_2$  (*weak solution*)
2. flow condition 3 according to formulas for normal shocks ( $M_{n_2} = M_2 \sin(\beta_2)$  and  $M_{n_3} = M_3 \sin(\beta_2 - \theta)$ )

⇒ complete description of flow and shock waves

(angle between upper wall and second shock:  $\Phi = \beta_2 - \theta$ )

# Roadmap - Oblique Shocks and Expansion Waves



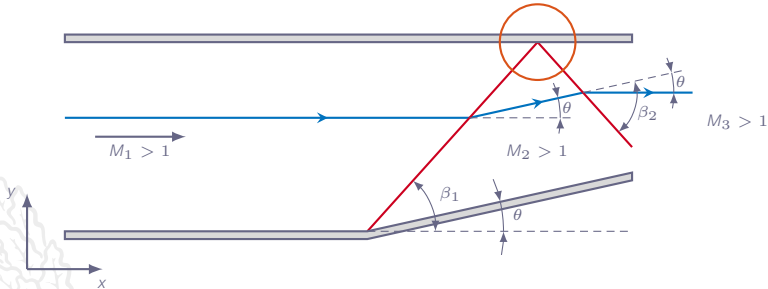
# Chapter 4.11

## Mach Reflection

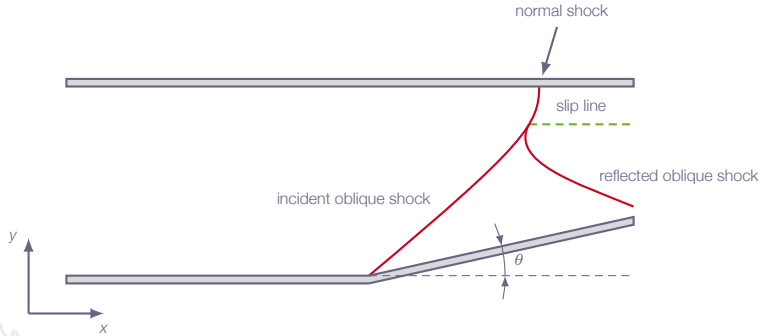


# Regular Shock Reflection

Regular reflection possible if both primary and reflected shocks are weak (see  $\theta$ - $\beta$ - $M$  relation)



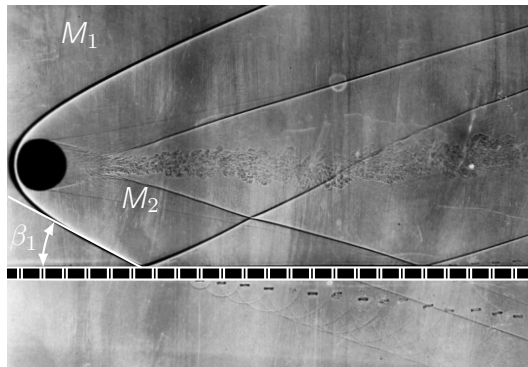
# Mach Reflection



Mach reflection:

appears when regular reflection is not possible  
more complex flow than for a regular reflection  
no analytic solution - numerical solution necessary

# Oblique Shocks and Mach Waves

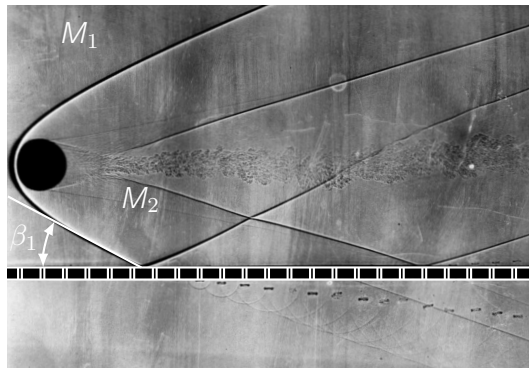


$$M_1 > M_2$$

$$M_2 > 1.0$$

$$\theta_1 = f(M_1, \beta_1), \quad M_2 = f(M_1, \theta_1, \beta_1)$$

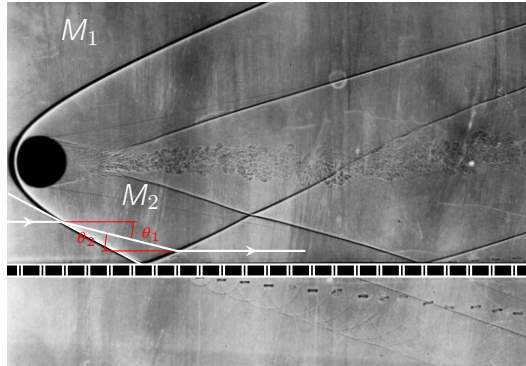
# Oblique Shocks and Mach Waves



$$\left. \begin{array}{l} \beta_1 = 28^\circ \\ M_1 = 3.1 \end{array} \right\} \Rightarrow \theta_1 \approx 11.2^\circ, \quad M_2 \approx 2.5$$



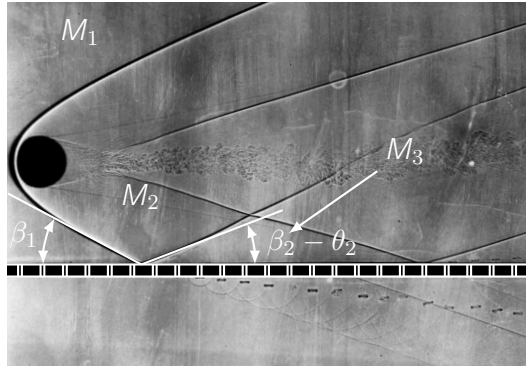
# Oblique Shocks and Mach Waves



$$\theta_1 = \theta_2$$



# Oblique Shocks and Mach Waves



$$M_1 > M_2 > M_3$$

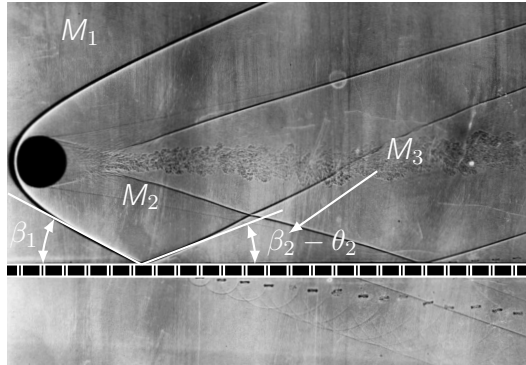
$$M_3 > 1.0$$

$$\beta_2 > \beta_1$$

$$\beta_2 = f(M_2, \theta_2), \quad M_3 = f(M_2, \theta_2, \beta_2)$$

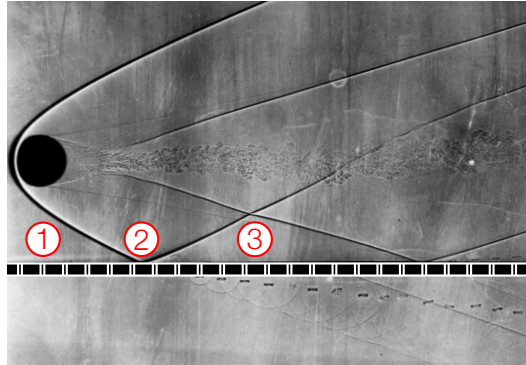
**Note!** Shock wave reflection at solid wall is **not** specular

# Oblique Shocks and Mach Waves



$$\left. \begin{array}{l} \theta_2 = 11.2^\circ \\ M_2 = 2.5 \end{array} \right\} \Rightarrow \beta_2 \approx 33^\circ, \quad M_3 \approx 2.0$$

# Oblique Shocks and Mach Waves



$$\frac{\rho_3}{\rho_1} = \frac{\rho_2}{\rho_1} \frac{\rho_3}{\rho_2} \approx 4.52$$

$$\frac{T_3}{T_1} = \frac{T_2}{T_1} \frac{T_3}{T_2} \approx 1.57$$

THE BERNOULLI-DOPPLER-LEIDENFROST-PELTZMAN-SAPIR-WHORE-DUNNING-KRUGER-STROOP EFFECT STATES THAT IF A SPEEDING FIRE TRUCK LIFTS OFF AND HURTLES TOWARD YOU ON A LAYER OF SUPERHEATED GAS, YOU'LL DIVE OUT OF THE WAY FASTER IF THE DRIVER SCREAMS "RED!" IN A *NON-TONAL* LANGUAGE THAT *HAS* A WORD FOR "FIREFIGHTER" THAN IF THEY SCREAM "GREEN!" IN A *TONAL* LANGUAGE WITH *NO* WORD FOR "FIREFIGHTER" WHICH YOU *THINK* YOU'RE FLUENT IN BUT *AREN'T*.

